# Physics-Informed Convolutional Neural Networks for Corruption Removal on Dynamical Systems

Daniel Kelshaw Department of Aeronautics Imperial College London djk21@imperial.ac.uk Luca Magri Department of Aeronautics Imperial College London, Alan Turing Institute 1.magri@imperial.ac.uk

## Abstract

Measurements on dynamical systems, experimental or otherwise, are often subjected to inaccuracies capable of introducing corruption; removal of which is a problem of fundamental importance in the physical sciences. In this work we propose physics-informed convolutional neural networks for stationary corruption removal, providing the means to extract physical solutions from data, given access to partial ground-truth observations at collocation points. We showcase the methodology for 2D incompressible Navier-Stokes equations in the chaotic-turbulent flow regime, demonstrating robustness to modality and magnitude of corruption.



Figure 1: Diagrammatic overview of the corruption removal task.

# 1 Introduction

In the physical sciences, corrupted data can arise for a number of reasons: faulty experimental sensors, or low-fidelity numerical results to name a few. The detection and removal of this corruption would allow for the extraction of the underlying true-state, ensuring that the governing equations are satisfied. In practice, verifying that given observations are characteristic of a dynamical system is more straightforward than determining the solution itself. For observations on dynamical systems we evaluate the residual, giving an indication of the fidelity of the observations.

Corruption removal typically considers noise-removal, and flow-reconstruction [1]. Noise-removal constitutes methods which remove small, stochastic variations from the true-state, with examples including: filtering methods [2], proper orthogonal decomposition [3, 4], and autoencoders [5]. Flow-reconstruction considers inference of a solution over the domain, given only partial information; a popular example being that of super-resolution. Convolutional neural networks (CNNs) are prominent in the domain of flow reconstruction due to their ability to exploit spatial correlations [6, 7].

In the absence of ground-truth labels, it is possible to impose prior knowledge by regularising predictions with respect to known governing equations [8]. Physics-informed CNNs have been used

Machine Learning and the Physical Sciences workshop, NeurIPS 2022.

for solving PDEs [9, 10], basic super-resolution [11], and surrogate modelling without access to labelled data [12]. This regularisation method differs from the standard physics-informed neural network (PINN) approach which strives to exploit the automatic-differentiation capabilities of neural networks to constrain gradients of the network [13].

Our work considers a form of physics-informed CNNs for flow reconstruction. We provide a method for removing arbitrary, stationary corruption from observations on dynamical systems, given access to only partial ground-truth information at select collocation points. We showcase corruption removal for the 2D incompressible Navier-Stokes equations in the chaotic-turbulent regime, highlighting the ability to accurately recover the true-state, independent of modality or magnitude of corruption.

## 2 Methodology

We consider corruption removal from observations on dynamical systems of the form

$$\partial_t \boldsymbol{u} - \mathcal{N}(\boldsymbol{u}; \lambda) = 0$$
 with  $\boldsymbol{u} = u(\boldsymbol{x}, t), \quad \boldsymbol{x} \in \Omega \subset \mathbb{R}^n, \quad t \in [0, T] \subset \mathbb{R}^+,$  (1)

where  $u: \Omega \times [0,T] \to \mathbb{R}^n$ ,  $\mathcal{N}$  is a sufficiently smooth differential operator, and  $\lambda$  are the physical parameters of the system. We define the residual of the system as the left-hand side of equation (1)

$$\mathcal{R}(\boldsymbol{u};\boldsymbol{\lambda}) \triangleq \partial_t \boldsymbol{u} - \mathcal{N}(\boldsymbol{u};\boldsymbol{\lambda}), \tag{2}$$

such that  $\mathcal{R}(\boldsymbol{u}; \lambda) = 0$  when  $u(\boldsymbol{x}, t)$  is a solution to the partial differential equation (PDE). We consider corrupted observations such that  $\mathcal{R}(\zeta(\boldsymbol{x}, t); \lambda) > 0$ , where  $\zeta(\boldsymbol{x}, t)$  is the linear combination of the true state  $u(\boldsymbol{x}, t)$  and the stationary corruption  $\phi(\boldsymbol{x})$ , precisely  $\zeta(\boldsymbol{x}, t) = u(\boldsymbol{x}, t) + \phi(\boldsymbol{x})$ .

Given the corrupted state, we aim to recover the underlying, physical solution to the governing equations: the true-state. Mathematically, we denote this process by the mapping

$$\eta_{\theta}: \zeta(\boldsymbol{\Omega}_{\boldsymbol{g}}, t) \to u(\boldsymbol{\Omega}_{\boldsymbol{g}}, t), \tag{3}$$

where the domain  $\Omega$  is discretised on a uniform, structured grid  $\Omega_g \subset \mathbb{R}^{N^n}$ . We further discretise the time domain, providing  $\mathcal{T} = \{t_i \in [0, T]\}_{i=0}^{N_t}$  for  $N_t$  samples in time. Approximating the mapping  $\eta_{\theta}$  as a CNN, we optimise weights of the network to minimise the loss

$$\mathcal{L}_{\theta} = \alpha \left( \mathcal{L}_{\partial \Omega} + \mathcal{L}_{\phi} \right) + \mathcal{L}_{\mathcal{R}},\tag{4}$$

where  $\alpha$  is a fixed, empirical weighting factor. We define each of the loss terms as

$$\mathcal{L}_{\mathcal{R}} = \frac{1}{|\mathcal{T}|} \sum_{t \in \mathcal{T}} \left\| \mathcal{R} \big( \eta_{\theta}(\zeta(\boldsymbol{\Omega}_{\boldsymbol{g}}, t)); \lambda \big) \right\|_{\boldsymbol{\Omega}_{\boldsymbol{g}}}^{2}, \tag{5}$$

$$\mathcal{L}_{\partial\Omega} = \frac{1}{|\mathcal{T}|} \sum_{t \in \mathcal{T}} \left\| \eta_{\theta}(\zeta(\partial \boldsymbol{\Omega}_{\boldsymbol{g}}, t)) - u(\partial \boldsymbol{\Omega}_{\boldsymbol{g}}, t) \right\|_{\partial \boldsymbol{\Omega}_{\boldsymbol{g}}}^{2}, \tag{6}$$

$$\mathcal{L}_{\phi} = \frac{1}{|\mathcal{T}|} \sum_{t \in \mathcal{T}} \left\| \partial_t \left[ \zeta(\boldsymbol{\Omega}_{\boldsymbol{g}}, t) - \eta_{\theta}(\zeta(\boldsymbol{\Omega}_{\boldsymbol{g}}, t)) \right] \right\|_{\boldsymbol{\Omega}_{\boldsymbol{g}}}^2, \tag{7}$$

where  $\partial \Omega_{g}$  denotes boundary points of the grid, and  $\|\cdot\|_{\Omega}$  is the  $\ell^{2}$ -norm over the given domain. We regularise network predictions with the residual-based loss  $\mathcal{L}_{\mathcal{R}}$ , imposing prior-knowledge of the system to promote network realisations which conform to the governing equations. The data-driven boundary loss  $\mathcal{L}_{\partial\Omega}$  serves to minimise the error between predictions and known measurements; when used in conjunction with the residual-loss, this effectively allows us to condition the underlying physics on the observations. For this case, we consider access to ground-truth measurements on the boundary  $\partial\Omega$ . Inclusion of the corruption-based loss  $\mathcal{L}_{\phi}$  embeds our assumption of stationary corruption, driving predictions away from the trivial solution,  $u(\Omega_{g}, t) = 0.^{\dagger}$ 

<sup>&</sup>lt;sup>†</sup>All code is available on GitHub: https://github.com/magriLab/PICR

## **3** Results

We consider the corruption removal for measurements on the Kolmogorov flow, an instance of the 2D incompressible Navier-Stokes equations. The flow is defined on the domain  $\Omega \in [0, 2\pi) \subset \mathbb{R}^2$ , imposing periodic boundary conditions on  $\partial\Omega$ , and subjected to stationary, sinusoidal forcing g(x). The chaotic-turbulent nature of the Kolmogorov flow provides a challenging case study, introducing complex non-linear dynamics and turbulent structures across a range of spatial frequencies [14]. The standard continuity and momentum equations are given as

$$\nabla \cdot \boldsymbol{u} = 0$$
  
$$\partial_t \boldsymbol{u} + \boldsymbol{u} \cdot \nabla \boldsymbol{u} = -\nabla p + \nu \Delta \boldsymbol{u} + g(\boldsymbol{x}),$$
(8)

where  $p, \nu$  denote the scalar pressure field and kinematic viscosity respectively. We take  $\nu = 1/34$  to ensure chaotic-turbulent dynamics, and prescribe the forcing model  $g(\boldsymbol{x}) = [\sin(4\boldsymbol{x}_2), 0]^{\top}$ .

## 3.1 Pseudospectral Discretisation

We provide a differentiable pseudospectral discretisation for the differential operator, defining operations on the spectral grid  $\hat{\Omega}_{k} \in \mathbb{Z}^{61 \times 61}$ , enabling backpropagation for the residual-based loss  $\mathcal{L}_{\mathcal{R}}$ . By eliminating the pressure term, our discretisation handles the continuity constraint implicitly, affording us the ability to neglect the constraint in the loss [15]. We produce a solution by time-integration of the dynamical system with the forward-Euler method, taking a time-step  $\Delta t = 0.005$  to ensure numerical stability according to the Courant-Friedrichs-Levy condition.

Due to the spectral discretisation, we evaluate the residual-based loss  $\mathcal{L}_{\mathcal{R}}$  in the Fourier domain on the grid  $\hat{\Omega}_k$ , offering two distinct advantages over the standard finite-difference approach: spectral accuracy, and implicit handling of periodic boundary conditions. We employ a form of time-windowed batching, considering the evaluation of the residual-based loss over  $\tau \geq 2$  consecutive time-steps.

### 3.2 Parameterised Corruption

We conduct experiments varying the frequency f, and relative magnitude  $\mathcal{M}$  of the corruption. For a scaling magnitude  $M = \mathcal{M} \cdot \max(u(\Omega_{g}, \cdot))$ , we define the corruption field as

$$\phi(\boldsymbol{x}; f, M) = \Gamma\left(20 + \sum_{i=1}^{2} \left[ (\boldsymbol{x}_i - \pi)^2 - 10\cos(f \cdot (\boldsymbol{x}_i - \pi)) \right]; M \right),$$
(9)

where  $\Gamma(\cdot; M)$  is a linear operator that re-scales the vector in the range [0, M], [16].



Figure 2: Corruption removal results for the Kolmogorov flow  $[f = 7, \mathcal{M} = 0.5]$ .

#### 3.3 Numerical Experiments

The Kolmogorov flow is solved in the Fourier domain, as per §3.1, before generating data on the grid  $\Omega_g \in \mathbb{R}^{64 \times 64}$ . A total of 1024 time-windows are used for training, with a further 256 for validation, taking  $\tau = 2$  in each instance. We employ a standard U-Net [17] architecture and train with the Adam optimizer, taking a learning rate of  $3 \times 10^{-4}$  for 5000 epochs, fixing computational resource for each experiment. We impose a fixed weighting of  $\alpha = 1 \times 10^3$  for the empirical weighting term.<sup>‡</sup>

In Figure 2 we demonstrate results for f = 7,  $\mathcal{M} = 0.5$ , showcasing the ability to separate the true-state from corruption. We note that the residual of the corruption is not equal to the sum of the residuals of it's constituents, highlighting the challenge of non-linearity. We further assess the robustness of the methodology to both modality and magnitude of corruption by running two studies:

(i) 
$$\mathcal{M} = 0.5, \quad f \in \{1, 3, 5, 7, 9\};$$
  
(ii)  $f = 3, \quad \mathcal{M} \in \{0.01, 0.1, 0.25, 0.5, 1.0\}.$ 

Results can be seen in Figure 3, where performance for all experiments is measured by the relative  $\ell^2$ -error between the predicted-state  $\eta_{\theta}(\zeta(\Omega_g, t))$ , and the true-state  $u(\Omega_g, t)$ . Results denote the mean over five repeats with random seeds, the standard deviation not exceeding  $8.169 \times 10^{-4}$ .



Figure 3: Comparison of relative  $\ell^2$ -error for parameterised corruption frequency and magnitude.

We find that performance is largely agnostic to both multi-modality and magnitude of corruption, with error decreasing slightly with increasing  $f, \mathcal{M}$ , as demonstrated in Figure 3. Quantitatively, we observe excellent robustness of the methodology, with relative  $\ell^2$ -error not exceeding  $2.284 \times 10^{-2}$ .

## 4 Conclusion

In this work we introduce a methodology for physics-informed removal of arbitrary, stationary corruption from observations on dynamical systems, given access to only partial ground-truth information. We demonstrate results for corrupted observations on the Kolmogorov flow, using ground-truth observations on the boundaries to condition the physics. Experiments indicate that performance is agnostic to both magnitude and modality of corruption, with relative  $\ell^2$ -error remaining consistently low, as demonstrated in Figure 3. Our results in Figure 2 qualitatively showcase the performance of the model, clearly separating the corrupted field into its component parts: true-state, and corruption. This work opens opportunities for accurate flow reconstruction from corrupted state measurements with access to partial ground-truth information, potentially allowing for correcting sensor measurements and low-fidelity simulation results - problems of fundamental importance in the physical sciences.

We note the limitations of our work: the assumption of stationary corruption, and the requirement of ground-truth observations to condition the physics; providing considerations for future work.

## Acknowledgments and Disclosure of Funding

D. Kelshaw. and L. Magri. acknowledge support from the UK EPSRC. L. Magri gratefully acknowledges financial support from the ERC Starting Grant PhyCo 949388.

<sup>&</sup>lt;sup>\*</sup>All experiments were run on a single NVIDIA Quadro RTX 8000.

## References

- [1] S. Brunton, B. Noack, and P. Koumoutsakos, "Machine learning for fluid mechanics," *Annual Review of Fluid Mechanics*, vol. 52, pp. 477–508, 1 2020.
- [2] N. Kumar and M. Nachamai, "Noise removal and filtering techniques used in medical images," Oriental journal of computer science and technology, vol. 10, pp. 103–113, 3 2017.
- [3] M. Raiola, S. Discetti, and A. Ianiro, "On piv random error minimization with optimal pod-based low-order reconstruction," *Experiments in Fluids*, vol. 56, 4 2015.
- [4] M. A. Mendez, M. Raiola, A. Masullo, S. Discetti, A. Ianiro, R. Theunissen, and J. M. Buchlin, "Pod-based background removal for particle image velocimetry," *Experimental Thermal and Fluid Science*, vol. 80, pp. 181–192, 1 2017.
- [5] P. Vincent, H. Larochelle, Y. Bengio, and P.-A. Manzagol, "Extracting and composing robust features with denoising autoencoders," in *Proceedings of the 25th International Conference on Machine Learning*, 2008, pp. 1096–1103.
- [6] C. Dong, C. C. Loy, K. He, and X. Tang, "Learning a deep convolutional network for image super-resolution," in *European Conference on Computer Vision*, 2014, pp. 184–199.
- [7] W. Shi, J. Caballero, F. Huszár, J. Totz, A. P. Aitken, R. Bishop, D. Rueckert, and Z. Wang, "Real-time single image and video super-resolution using an efficient sub-pixel convolutional neural network," *Computer Vision and Pattern Recognition*, 9 2016.
- [8] I. E. Lagaris, A. Likas, and D. I. Fotiadis, "Artificial neural networks for solving ordinary and partial differential equations," *IEEE Transactions on Neural Networks*, vol. 9, pp. 987–1000, 1998.
- [9] H. Gao, L. Sun, and J. X. Wang, "Phygeonet: Physics-informed geometry-adaptive convolutional neural networks for solving parameterized steady-state pdes on irregular domain," *Journal of Computational Physics*, vol. 428, 3 2021.
- [10] P. Ren, C. Rao, Y. Liu, J. X. Wang, and H. Sun, "Phycrnet: Physics-informed convolutionalrecurrent network for solving spatiotemporal pdes," *Computer Methods in Applied Mechanics and Engineering*, vol. 389, 2 2022.
- [11] H. Gao, L. Sun, and J.-X. Wang, "Super-resolution and denoising of fluid flow using physicsinformed convolutional neural networks without high-resolution labels," *Physics of Fluids*, vol. 33, p. 073603, 7 2021.
- [12] X. Zhao, Z. Gong, Y. Zhang, W. Yao, and X. Chen, "Physics-informed convolutional neural networks for temperature field prediction of heat source layout without labeled data," *Preprint*, 9 2021. [Online]. Available: http://arxiv.org/abs/2109.12482
- [13] M. Raissi, P. Perdikaris, and G. E. Karniadakis, "Physics-informed neural networks: A deep learning framework for solving forward and inverse problems involving nonlinear partial differential equations," *Journal of Computational Physics*, vol. 378, pp. 686–707, 2 2019.
- [14] E. D. Fylladitakis, "Kolmogorov flow: Seven decades of history," Journal of Applied Mathematics and Physics, vol. 06, pp. 2227–2263, 2018.
- [15] C. Canuto, M. Y. Hussaini, A. Quarteroni, and T. A. Zang, *Spectral Methods in Fluid Dynamics*. Springer Berlin Heidelberg, 1988.
- [16] L. A. Rastrigin, "Systems of extremal control," Nauka, 1974.
- [17] O. Ronneberger, P. Fischer, and T. Brox, "U-net: Convolutional networks for biomedical image segmentation," in *Medical Image Computing and Computer-Assisted Intervention*, 2015, pp. 234–241.

## Checklist

- 1. For all authors...
  - (a) Do the main claims made in the abstract and introduction accurately reflect the paper's contributions and scope?

[Yes] We claim to provide a methodology which removes arbitrary, stationary corruption from observations on dynamical systems. We quantitatively demonstrate this in Figure 3, and qualitatively in Figure 2.

- (b) Did you describe the limitations of your work?
   [Yes] We note limitations of our work in the conclusion, §4. Namely, the assumption of stationary corruption, and the dependence on ground-truth observations at collocation points.
- (c) Did you discuss any potential negative societal impacts of your work?[No] We do not for-see any negative societal impacts of the work.
- (d) Have you read the ethics review guidelines and ensured that your paper conforms to them? [Yes]
- 2. If you are including theoretical results...
  - (a) Did you state the full set of assumptions of all theoretical results? [N/A]
  - (b) Did you include complete proofs of all theoretical results? [N/A]
- 3. If you ran experiments...
  - (a) Did you include the code, data, and instructions needed to reproduce the main experimental results (either in the supplemental material or as a URL)?
     [Yes] We include a link to the code in the footprint of §2.
  - (b) Did you specify all the training details (e.g., data splits, hyperparameters, how they were chosen)?

[Yes] We provide all necessary training details in §3, including: architecture, hyperparameters, amount of data, physical parameters for simulation results. Further details are available on the link to the repository.

(c) Did you report error bars (e.g., with respect to the random seed after running experiments multiple times)?

[Yes] We state an upper-bound of the standard deviation for repeats across five random seed in §3. The value is so small it does not warrant putting error bars in Figure 3.

(d) Did you include the total amount of compute and the type of resources used (e.g., type of GPUs, internal cluster, or cloud provider)?

## [Yes] We provide a footnote in §3.3 which states the computational resource used.

- 4. If you are using existing assets (e.g., code, data, models) or curating/releasing new assets...
  - (a) If your work uses existing assets, did you cite the creators? [N/A]
  - (b) Did you mention the license of the assets? [N/A]
  - (c) Did you include any new assets either in the supplemental material or as a URL?
     [Yes] We provide a differentiable pseudospectral spatial discretisation for the Kolmogorov flow, a link to the repository is available in §2.
  - (d) Did you discuss whether and how consent was obtained from people whose data you're using/curating? [N/A]
  - (e) Did you discuss whether the data you are using/curating contains personally identifiable information or offensive content? [N/A]
- 5. If you used crowdsourcing or conducted research with human subjects...
  - (a) Did you include the full text of instructions given to participants and screenshots, if applicable? [N/A]
  - (b) Did you describe any potential participant risks, with links to Institutional Review Board (IRB) approvals, if applicable? [N/A]
  - (c) Did you include the estimated hourly wage paid to participants and the total amount spent on participant compensation? [N/A]