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# A machine learning approach to duality in statistical physics

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## Abstract

The notion of *duality* – that a given physical system can have two different mathematical descriptions – is a key idea in modern theoretical physics. Establishing a duality in lattice statistical mechanics models requires the construction of a dual Hamiltonian and a map from the original to the dual observables. By using neural networks to parameterize these maps and introducing a loss function that penalises the difference between correlation functions in original and dual models, we formulate the process of duality discovery as an optimization problem. We numerically solve this problem and show that our framework can rediscover the celebrated Kramers-Wannier duality for the 2d Ising model, numerically reconstructing the known mapping of temperatures.\* We discuss future directions and prospects for discovering new dualities within this framework.

## 1 Background

A key concept in physics is *duality*, i.e. the idea that the same physical system can have two different mathematical descriptions. Duality sits at the heart of modern theoretical physics. In this work we seek to formalize the notion of duality in statistical physics in a manner that allows modern machine learning techniques to be used to systematically search for dualities.

**Background on duality:** Consider a statistical physics model with microstates  $\sigma$  and Hamiltonian functional  $H[\beta, \sigma]$ , where  $\beta$  are macroparameters such as the temperature. The model is determined by its partition function  $Z = \sum_{\sigma} e^{-H[\beta, \sigma]}$ . However, in nature we often have access to sets of expectation values of observables  $O_{\alpha}(\sigma)$  (some real-valued functions of the microstates, e.g. correlation functions, with  $\alpha$  being an arbitrary label)

$$\langle O_{\alpha}(\sigma) \rangle_H = \frac{1}{Z} \sum_{\sigma} O_{\alpha}(\sigma) \exp(-H[\beta, \sigma]). \quad (1)$$

It is a profound physical fact that occasionally there are alternative representations of these sets of correlation functions (see e.g. [1; 2; 3; 4; 5; 6; 7] for influential examples). That is, there exists another set of microstates  $\tilde{\sigma}$ , another Hamiltonian  $\tilde{H}[\tilde{\beta}, \tilde{\sigma}]$  and for each observable  $O_{\alpha}(\sigma)$  a dual

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\*Our implementation is publicly available at [https://github.com/pg2455/duality\\_exploration](https://github.com/pg2455/duality_exploration)

observable  $\tilde{O}_\alpha(\tilde{\sigma}_i)$  such that

$$\langle O_\alpha(\sigma) \rangle_H = \langle \tilde{O}_\alpha(\tilde{\sigma}) \rangle_{\tilde{H}}. \quad (2)$$

When this happens, we have a *duality* –the same physical system has at least two distinct mathematical descriptions, which may be useful for different reasons.

A prototypical example of such a duality is Kramers-Wannier duality for the 2d Ising model [2]. The 2d Ising model† consists of spins  $\sigma_i = \pm 1$  living on the sites of a square lattice at temperature  $\beta^{-1}$ , with Hamiltonian that sums over neighbouring spins  $\langle ij \rangle$

$$H[\beta, \sigma] = -\beta \sum_{\langle ij \rangle} \sigma_i \sigma_j. \quad (3)$$

It is a remarkable fact that the model described by (3) is precisely equivalent to a different 2d Ising model with spins  $\tilde{\sigma}_i = \pm 1$  living on the *dual* lattice, with a dual Hamiltonian of the same functional form  $\tilde{H}[\tilde{\beta}, \tilde{\sigma}] = H[\beta, \tilde{\sigma}] = -\tilde{\beta} \sum_{\langle ij \rangle} \tilde{\sigma}_i \tilde{\sigma}_j$ , but with  $\tilde{\beta}$  satisfying  $\sinh(2\beta) \sinh(2\tilde{\beta}) = 1$ . The fact that the functional form of the Hamiltonian is the same is exceptional, and in this case one can call the duality a *self*-duality.

Importantly, all observables constructed from the  $\sigma_i$  can be mapped to observables of the  $\tilde{\sigma}_i$ . Consider for instance two neighbouring spins  $\sigma_i$  and  $\sigma_j$ . We can build an observable  $O_{ij} = \sigma_i \sigma_j$ , which we call a *link product*. Then the KW duality implies that

$$\langle O_{ij} \dots \rangle_H = \langle \tilde{O}_{ij}(\tilde{\sigma}) \dots \rangle_{\tilde{H}}, \quad \tilde{O}_{ij}(\tilde{\sigma}) = e^{-2\tilde{\beta}\tilde{\sigma}_{i^*}\tilde{\sigma}_{j^*}} \quad (4)$$

where the notation  $\tilde{\sigma}_{i^*}$  refers to sites on the dual lattice such that the link connecting sites  $i^*$  and  $j^*$  intersects the link connecting  $i$  and  $j$ . The  $\dots$  indicate that this is an operator equation which holds for arbitrary insertions of operators and thus can be used to construct any expectation value of an even number of the  $\sigma_i$ . Appropriate products of the link products determine all correlation functions.‡

In this work we tackle the problem of finding dualities using machine learning. In particular, starting from the original model  $(H, O_{ij})$  we formulate an optimization problem whose solution recovers *both*  $(H, O_{ij})$  as well as  $(\tilde{H}, \tilde{O}_{ij})$ . This constitutes an automated discovery of a duality.

**Previous work:** The problem of learning the parameters in a Hamiltonian from data is precisely that of training a Boltzmann machine, and has a very long history. Our case differs from the classical situation in that we are simultaneously learning a mapping of observables. Other work on dualities involving machine learning includes [9], [10], but none is aimed at recovering the full dual descriptions as we do here.

## 2 Methodology

We now explain how, starting from the Hamiltonian  $H[\beta, \sigma]$  of some statistical model on a lattice, we can learn candidates  $\tilde{H}[\tilde{\beta}, \tilde{\sigma}]$  for dual models as well as a dictionary between original and dual observables. This includes learning the fact that the dual model is defined on the dual lattice.

**Framework and loss function:** We assume that  $\tilde{H}$  can be written in terms of local couplings of spins:

$$\tilde{H}[\tilde{\beta}, \tilde{\sigma}] = -\tilde{\beta} \sum_{\langle ij \rangle} \tilde{\sigma}_i \tilde{\sigma}_j - \tilde{\beta}_4 \sum_{\langle ijkl \rangle} \tilde{\sigma}_i \tilde{\sigma}_j \tilde{\sigma}_k \tilde{\sigma}_l - \dots \quad (5)$$

where the couplings  $\tilde{\beta}_a$  are parameters to be learned. We would like to find dual representations of the link products  $O_{ij}$  we described for the Ising model. We assume that the link product in the

† See e.g. [8] for a textbook treatment.

‡ Due to a  $\mathbb{Z}_2$  symmetry the expectation value of a moment of an odd number of spins *formally* vanishes in a finite model, though as usual in the symmetry spontaneously broken phase this might not be observed in a simulation that uses local updates. We also note that the relation is modified if we consider precisely the same two-spin operator *twice*, i.e.  $(\sigma_i \sigma_j)^2 = 1$ , when a careful derivation of the duality shows that the right-hand side must be modified and is also identically 1.

original model is mapped to *some* functions of *nearby* link products in the dual model, more precisely

$$\tilde{O}_{ij}(\tilde{\sigma}) = G(\{\tilde{\sigma}_k \tilde{\sigma}_l\}) \quad (6)$$

where  $\{\tilde{\sigma}_k \tilde{\sigma}_l\}$  is a set of link products such as the one shown in Figure 1.

$G$  is designed to be sufficiently flexible to recover models on lattices related in various ways to the original one. Note that a choice must be made about how to relate the assignment of link products neighbouring a horizontal link to the assignment of link products neighbouring a vertical link, as multiple choices are consistent with rotational invariance. In Figure 1 we display the choice used, which relates them by a rotation composed with a reflection. As we will see later, this choice is important for recovering the geometry of the dual lattice.

We now construct a loss function  $\mathcal{L}$  that is minimized when all correlation functions of  $O_{ij}$  and  $\tilde{O}_{ij}$  agree on the two sides of the duality. This is similar to the matching of moments of two distributions, which is a standard problem, and for which one can construct general kernels that are minimized only when all of the moments of two distributions agree (see e.g. [11]). Unfortunately, in the present case we cannot use kernels because of one conceptual and one technical problem: 1) certain moments need not be matched, as per Footnote 2, and 2) no notion of locality is embedded in standard moment matching (in the present case, correlation functions of faraway spins carry little information).

Instead we explicitly match *features* – i.e. moments of a small number of nearby link products, as shown in Figure 2 – which we then spatially average over the lattice. Denoting these features as  $\phi^a$  with  $a$  running over features, we then construct the loss

$$\mathcal{L}(G, \tilde{H}) = \sum_a \ell^a \ell^a \quad \ell^a = \langle \phi^a[\sigma_i] \rangle_H - \langle \phi^a[G(\tilde{\sigma}_i)] \rangle_{\tilde{H}} \quad (7)$$

$\ell^a$  can be thought of as a vector in feature space indicating how far apart the two theories are.

It is clear that this loss can be minimized in two scenarios: (a)  $\tilde{H} = H$  and  $G(\tilde{\sigma}_i \tilde{\sigma}_j) = \tilde{\sigma}_i \tilde{\sigma}_j$ , i.e., the original model is rediscovered, or (b)  $\tilde{H} \neq H$  and  $G(\tilde{\sigma}_i \tilde{\sigma}_j) \neq \tilde{\sigma}_i \tilde{\sigma}_j$ , representing a nontrivial dual model where (selected) moments nevertheless perfectly match those of the original model.

**Optimization:** We now need to solve the following optimization problem:

$$G^*, \tilde{H}^* = \arg \min_{G, \tilde{H}} \mathcal{L}(G, \tilde{H}) \quad (8)$$

$G$  is represented by a neural network with parameters  $\theta$ ,  $G = G_\theta$ .

Algorithm 1 outlines the procedure for optimization. Given a trial set of parameters  $\theta$  and couplings for the dual Hamiltonian  $\tilde{\beta}_a$ , we simultaneously perform Markov Chain Monte Carlo (MCMC) sampling from the original and dual Hamiltonians using a standard Metropolis algorithm to obtain spin configurations  $\sigma_i$  and  $\tilde{\sigma}_i$  drawn from the appropriate distributions respectively. We can then evaluate the expectation values in (7), and compute the loss  $\mathcal{L}$ .

To minimize it we also need to compute gradients  $\partial_\theta \mathcal{L}$  and  $\partial_{\tilde{\beta}_a} \mathcal{L}$ . For  $\theta$  this can be done straightforwardly using conventional automatic differentiation techniques. For the  $\tilde{\beta}_a$  we cannot backpropagate through a stochastic sampler, but explicit differentiation shows that we can relate the gradients to expectation values that can be evaluated through MCMC sampling from the dual Hamiltonian (see

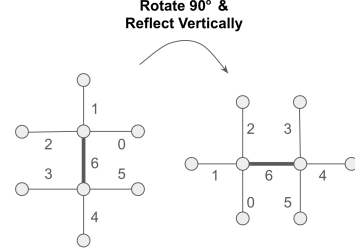


Figure 1: We parametrize  $G$  as a neural network that takes neighboring links of a given link (in this case # 6) as its input. The assignment on horizontal links is related to that on vertical ones by a rotation and reflection.



Figure 2: Examples of three features showing link products considered.

Appendix C for derivation), e.g. <sup>§</sup>

$$\partial_{\tilde{\beta}} \mathcal{L} = 2 \left\langle \sum_a \ell^a \left( \sum_{\langle ij \rangle} \langle \sigma_i \sigma_j \rangle_{\tilde{H}} - \sum_{\langle ij \rangle} \tilde{\sigma}_i \tilde{\sigma}_j \right) \phi^a [G_\theta(\tilde{\sigma})] \right\rangle_{\tilde{H}}. \quad (9)$$

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**Algorithm 1** Machine learning for finding statistical mechanics duality

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- 1: **Inputs:**  $\beta, \eta$  (learning rate),  $N$  (number of samples)
  - 2: **Initialize:**  $\tilde{\beta}_0 \in \mathbb{R}, \theta \in \mathbb{R}^d$
  - 3: **for** each epoch  $t = 1, 2, \dots, T$  **do**
  - 4:   Draw  $N$  samples  $\{\sigma_i\}_{i=1}^N \sim p(\sigma|\beta)$
  - 5:   Draw  $N$  samples  $\{\tilde{\sigma}_i\}_{i=1}^N \sim p(\tilde{\sigma}|\tilde{\beta})$  where  $\tilde{\beta} \neq \beta$
  - 6:   Compute the loss  $\mathcal{L} = \frac{1}{N} \sum_{i=1}^N \mathcal{L}(\sigma_i, G_\theta(\tilde{\sigma}_i))$
  - 7:   Compute the gradients  $\partial_{\tilde{\beta}} \mathcal{L}$  and  $\partial_\theta \mathcal{L}$
  - 8:   Update the parameters:
 
$$\begin{aligned} \tilde{\beta}_{t+1} &\leftarrow \tilde{\beta}_t - \eta \partial_{\tilde{\beta}} \mathcal{L} \\ \theta_{t+1} &\leftarrow \theta_t - \eta \partial_\theta \mathcal{L} \end{aligned}$$
  - 9:   **if**  $\mathcal{L}$  has not improved for the last  $X$  epochs **then**
  - 10:     **Stop the optimization**
  - 11:   **end if**
  - 12: **end for**
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### 3 Experiments

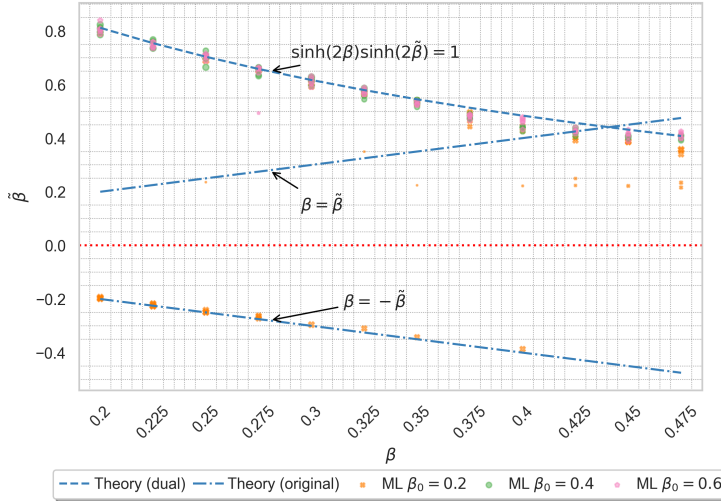


Figure 3: Final  $\tilde{\beta}$  as found by the deep learning framework closely matches that of the theoretical results. Points are scaled by the negative logarithm of the best loss such that the **size of the points is inversely proportional to the loss**. We cap the minimum size so that smaller points are visible. The loss is a minimum along two fronts, i.e., original frame  $\beta = \pm \tilde{\beta}$  and the dual frame along the lines  $\sinh(2\beta) \sinh(2\tilde{\beta}) = 1$ .

<sup>§</sup>We note that this evaluation is computationally expensive, as each gradient step requires us to equilibrate an MCMC chain. For training conventional Boltzmann machines one can use more efficient approaches such as contrastive divergence [12]. Due to the presence of the mapping  $G$ , we are not aware of a similarly efficient algorithm in our case, and indeed all likelihood-based approaches seem conceptually difficult.

In this section, we describe some simple experiments using the above machinery. We take our original Hamiltonian  $H$  to be that of the 2d Ising model (3), and we take the dual Hamiltonian  $\tilde{H}$  in (5) to have only one non-zero parameter  $\tilde{\beta}$  (and so  $\tilde{\beta}_4 = 0$ , etc.).

**Neural Network architecture for  $G_\theta$ :** For a given link product in the dual frame we assemble the 7 nearby links shown in Fig. 1 into a 7-dimensional vector  $\mathbf{f}_{\langle ij \rangle} \in (\mathbb{Z}_2)^7$ , where each element of the vector is the product of the two spins living on the two ends of the link. We consider a simplistic neural network acting on this input, with parameters formed by  $\theta_1 \in \mathbb{R}^7$ , and

scalars  $\theta_2$  and  $\theta_3$ . We opt for hard attention using Gumbel-Softmax [13] so that only a few of the seven nearby links are utilized in the prediction task. Thus, the mapping is defined by,

$$G_\theta(\mathbf{f}_{\langle ij \rangle}) = \theta_2 \cdot \text{Gumbel-Softmax}(\theta_1)^T \mathbf{f}_{\langle ij \rangle} + \theta_3 \quad (10)$$

As the elements of  $\mathbf{f}_{\langle ij \rangle}$  are  $\pm 1$ , a very simple network provides a very expressive function.

**Rediscovery of the 2d Ising duality.** In Figure 3, we show the result of deploying the above machinery on different model values of  $\beta$  on an  $8 \times 8$  lattice with periodic boundary conditions. For each value of the input  $\beta$ , we ran a total of 15 optimizations, five from each of three initializations of  $\tilde{\beta}$ , i.e.,  $\tilde{\beta}_0 = 0.2$ ,  $\tilde{\beta}_0 = 0.4$  and  $\tilde{\beta}_0 = 0.6$ . Due to the randomness involved in MCMC sampling, each seed is expected to be an independent run.

We record the value of  $\tilde{\beta}$  obtained. There are three branches of solutions: the original model  $\tilde{\beta} = \beta$ , the dual model  $\sinh(2\beta) \sinh(2\tilde{\beta}) = 1$ , and an antiferromagnetic analogue of the original model  $\tilde{\beta} = -\beta$ . The latter is equivalent to the original frame, and is obtained by making the change of variables  $\sigma_i \rightarrow -\sigma_i$  on every other site, thus flipping the sign of  $\beta \rightarrow -\beta$ . Note that the existence of the dual branch of solutions can be viewed as a numerical “rediscovery” of the KW duality line

$$\sinh(2\beta) \sinh(2\tilde{\beta}) = 1 \quad (11)$$

Further details on the experiments (including an exploration on how they depend on the system size) are shown in the Supplementary Material.

It is interesting to ask how the model recovers the structure of the dual lattice, as well as the dual observables. The attention mechanism used encourages the model to use only a single link of the input, and for the runs that find the dual temperature this ends up using the links numbered either 2 or 5 instead of the original 6 in Figure 1. As we show in an example in Figure 4, this is equivalent to finding the dual lattice from the original. Note that here it is important that we relate horizontal to vertical links by the composition of a rotation *and* reflection as shown in Figure 1; other choices will not result in the possibility of finding the dual lattice, and indeed in our experiments they do not find a duality. The optimized values of  $G_\theta$  closely match theoretical results  $\tilde{O}_{ij}(\tilde{\sigma}) = e^{-2\tilde{\beta}\tilde{\sigma}_{i*}\tilde{\sigma}_{j*}}$ , as shown in more detail in the supplementary material.

In this approach the 1-1 mapping of  $\beta$  to  $\tilde{\beta}$  is only found numerically; one could possibly supplement this numerical determination with symbolic regression [14] to obtain an analytic formula such as (11), but in more complicated examples of the duality we do not expect there to necessarily exist a simple analytic formula and thus have not explored this.

## 4 Conclusions

Above we have explained how the process of finding dualities can be automated, demonstrating the mechanism by “rediscovering” the well-known Kramers-Wannier duality of the 2d Ising model. This is only a proof of principle, and much work remains to be done.

For example, as discussed in Section 3, at present we match a number of features which are constructed by hand. It would be ideal to find a kernel that allows matching of all the required moments while simultaneously giving lower weight to those involving faraway spins. On the operational side, it would be helpful to have a more efficient way of training; contrastive divergence fails here as there appears to be no simple way to map the likelihood of a single spin configuration across the duality.

On the physics side, we hope to use such techniques to find entirely new dualities. One concrete direction is to search for Kramers-Wannier duals of deformed Ising models, where extra spin-spin couplings have been added to the action: while some results exist for specific models [15; 16; 17], we are not aware of a completely general approach. In the setup above our preliminary experiments show that adding new couplings generically significantly hurts performance, and a more robust network

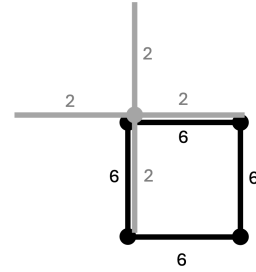


Figure 4: Emergence of dual lattice: e.g. if four original links (marked by 6) form a square, the corresponding four links that are referenced by the neighbour mapping (marked by 2) in Figure 1 form a cross, as expected for the dual lattice.

architecture and training dynamics is desirable. We are currently investigating approaches which explicitly use more of the known symmetry structure of Kramers-Wannier duality. A less concrete but far more exciting direction would be if one could use the approach to find entirely new dualities, unconnected to any existing ones. We hope to return to this in the future.

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## A Supplementary material

We provide some further details on our experimental results. A rough measure of the uncertainty of our results can be obtained from the spread of the learned  $\tilde{\beta}$  points in Figure 3, which grows as we approach the phase transition at  $\beta_c \approx 0.44$ . Interestingly, the method does not perform reliably for  $\beta > \beta_c$ , when the original frame is in the symmetry-broken phase. This is somewhat reminiscent of known difficulties in learning parameters of Hamiltonians at high  $\beta$  (see e.g. Appendix B of [18]) and deserves further study.

Further, from Figure 3 we observe that the initialization of  $\beta$  has an impact on the frame that is chosen. We find that  $\beta_0 = 0.4$  almost always drifted away from the original frame as indicated by the absence of green points on the  $\beta = -\tilde{\beta}$  line. Note for some  $\beta$  (e.g.,  $\beta = 0.35$ ), we see some suboptimal runs resulting in the final  $\beta$  far away from the dual or original frame. These have high loss and can be easily identified as failed runs.

Figure 5 shows runs from  $\beta = 0.2$  grouped by  $\beta_0$  and  $\tilde{\beta}^*$ , illustrating how the training progresses under different scenarios. For the seeds where either the dual or original frame is recovered, the loss goes to 0. Further, we track the entropy of Gumbel-Softmax( $\theta_1$ ) to assess how the algorithm is weighing each feature. A value of 0 corresponds to a strong preference for one out of the seven input links.

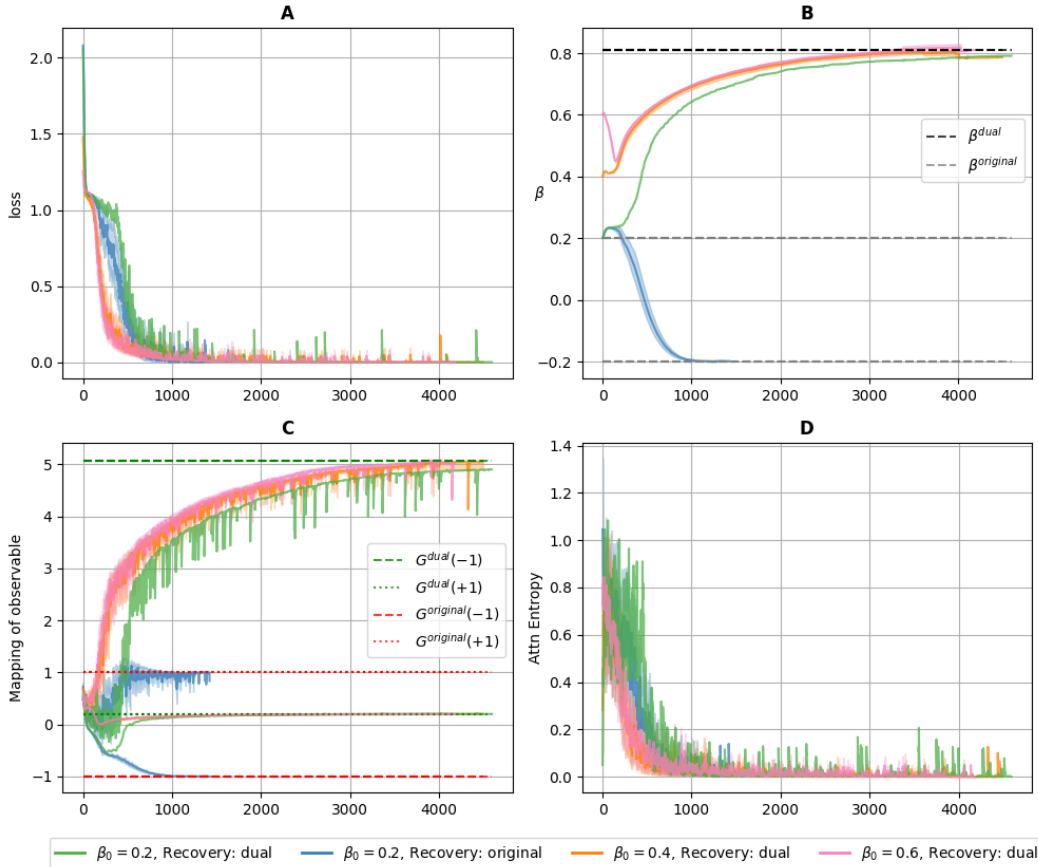


Figure 5: Training progress for runs from  $\beta = 0.2$ , grouped by  $\beta_0$  and  $\tilde{\beta}^*$  to showcase the trajectory of various metrics. We show exponentially smoothed moving average of the following metrics: (A) Loss, (B)  $\beta$ , (C) Mapping of observables, (D) Entropy of Gumbel-Softmax( $\theta_1$ ). For (B) and (C) we denote theoretically expected values in original and dual frames by the dashed lines.

Finally, in Figure 6 we see which link is selected as a map to the observable, using the numbering in Figure 1; links 2 and 5 correspond to a mapping to the dual lattice (and are found when we recover a map to the dual frame) and link 6 corresponds to recovering the original link (and are found when we



recover the original frame). The small amounts of other dimensions correspond to failed runs which generically have a higher loss.

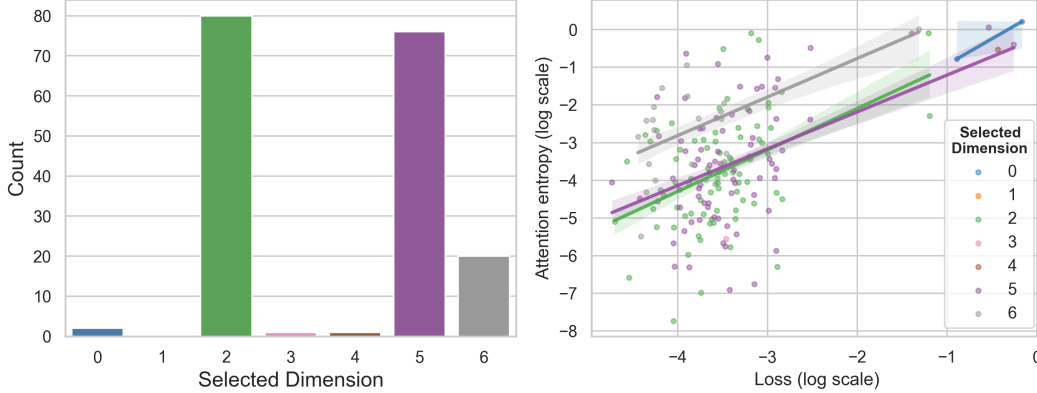


Figure 6: (left) Corresponding to Figure 3, we plot the frequency with which the dimension of the feature vector  $\mathbf{f}_{(ij)}$  was selected. Note that 2 and 5 pertain to the dual frame, while 6 relates to the discovery of the original frame. (right) Loss values (log scale) and attention entropy (log scale) are positively correlated such that lower loss increasingly prefers a single dimension of feature vectors. Note that both the loss and attention entropy are very low on the three features corresponding to the dual and original frames, as expected for a successful run.

## B Neural Network Training

Our models are all implemented in PyTorch [19]. We used the Adam [20] optimizer with the learning rate of 0.005 for  $\beta$  and 0.01 for  $\theta$ . [NI: Note I changed this so that it is correct – the effective learning rate from the algorithm was off by a factor of 2 for  $\beta$  due to the error previously]. Moreover, we used the early stopping criterion to stop the training if the loss didn't improve over 200 epochs. We ran the sampler in each experiment to generate 1000 samples for the lattice. We ran the training for a maximum of 25000 epochs, and our runs took about 1-3 hours each. Our experiments are run on the lattice size of  $8 \times 8$ .

## C Derivation

Here we derive (9). We seek to compute the gradients with respect to  $\tilde{\beta}$  of  $\mathcal{L}$  defined in (7). Recall that for any function of spins  $\mathcal{O}[\tilde{\sigma}]$  we have

$$\langle \mathcal{O} \rangle_{\tilde{H}} \equiv \frac{1}{Z(\tilde{\beta})} \sum_{\{\tilde{\sigma}_i\}} \mathcal{O}[\tilde{\sigma}] e^{(\sum_{\langle ij \rangle} \tilde{\beta} \tilde{\sigma}_i \tilde{\sigma}_j)} \quad Z(\tilde{\beta}) \equiv \sum_{\{\tilde{\sigma}_i\}} e^{(\sum_{\langle ij \rangle} \tilde{\beta} \tilde{\sigma}_i \tilde{\sigma}_j)} \quad (12)$$

where the sum over  $\{\sigma_i\}$  runs over all spin configurations. Now we have

$$\partial_{\tilde{\beta}} \mathcal{L} = -2 \sum_a \ell^a \partial_{\tilde{\beta}} \langle \phi^a [G(\tilde{\sigma}_i)] \rangle_{\tilde{H}}, \quad (13)$$

where we have used the definition of  $\ell^a$  in (7). From (12) the gradient of any observable with respect to  $\tilde{\beta}$  is

$$\partial_{\tilde{\beta}} \langle \mathcal{O} \rangle_{\tilde{H}} = -\langle \mathcal{O} \rangle_{\tilde{H}} \sum_{\langle ij \rangle} \langle \tilde{\sigma}_i \tilde{\sigma}_j \rangle_{\tilde{H}} + \sum_{\langle ij \rangle} \langle \tilde{\sigma}_i \tilde{\sigma}_j \mathcal{O} \rangle_{\tilde{H}} \quad (14)$$

where the first term comes from differentiating  $Z(\tilde{\beta})$  and the second from differentiating inside the Boltzmann measure weighting each configuration in (12). Using this expression to evaluate (14) for  $\mathcal{O} = \phi^a [G(\tilde{\sigma}_i)]$  we find (9).

## D Scaling results

Figure 7 shows the fraction of instances in which either  $\tilde{\beta}$ ,  $\beta$ , or  $-\beta$  were successfully recovered. We observe that as the lattice size increases to  $10 \times 10$  and  $12 \times 12$ , the recovery rate improves.

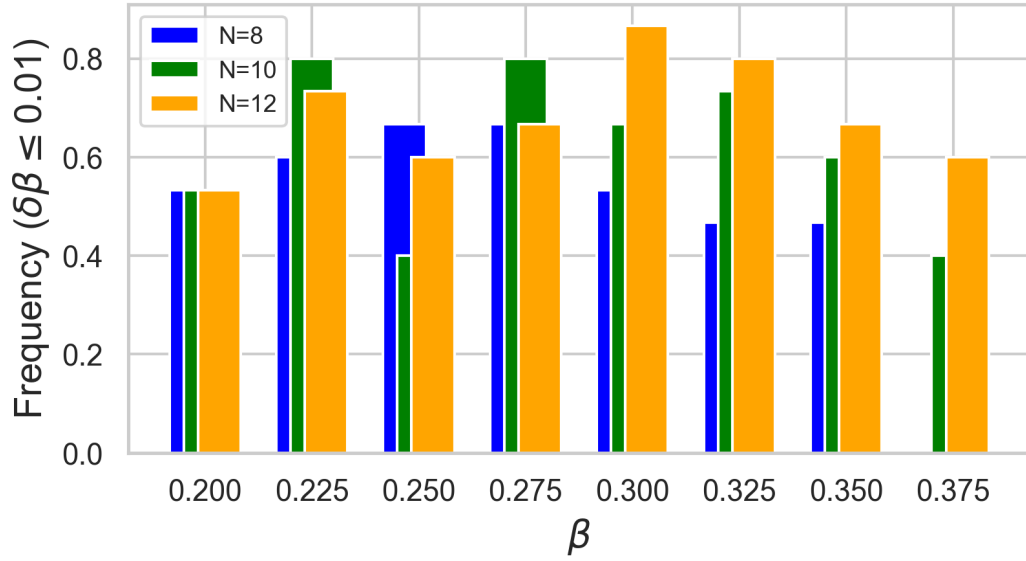


Figure 7: Fraction of times relevant beta was recovered within the tolerance of 0.01 as the lattice size is increased.

## E Broader Impact

Our work, on further scaling and addressing the current limitations, hold significant potential to advance knowledge in statistical physics. At this stage, we do not anticipate any negative societal impacts.

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